Phase transition in the Ising model on a random 2D lattice

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joint work with Joonas Turunen

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Linxiao Chen

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Introduction

Linxiao Chen

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Input: • A finite graph *G* embedded in \mathbb{R}^2 . • A coupling constant $\nu > 0$.

The *nearest-neighbor Ising model* on the *faces* of *G* is a random assignment of the numbers +1 and -1 (spins) to the faces of *G* according to the probability distribution such that for all $\sigma \equiv (\sigma_x)_{x \in F(G)} \in \{+1, -1\}^{F(G)}$,

 $\mathbb{P}(\sigma) \propto \nu^{\#\{x \sim y: \, \sigma_x = \sigma_y\}} =: \nu^{\mathcal{E}(G,\sigma)}$



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The coupling constant is related to the *physical parameters* by $\nu = \exp\left(\frac{2\beta}{k_{B}T}\right)$.

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 \sim large ν : "low temperature", small ν : "high temperature".

The usually studied case:

- $\nu > 1$ (ferromagnetic).
- *G* is a subgraph of a regular 2D lattice $(\triangle, \Box \text{ or } \bigcirc)$.

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• Demonstrate the existence of *phase transition* in a mathematically tractable model.

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Why would this extremely idealized model be related to the real physics?
 → <u>Universality</u>: the phenomenon that <u>some macroscopic observables</u> of the system <u>at or near criticality</u> is independent from microscopic details of the system.
 → Distinction between <u>non-universal observables</u> and <u>universal observables</u>.

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• Non-analyticity at $\nu = \nu_c$ of the *free energy density* when the system size $\rightarrow \infty$.

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: $\mathbb{E}[\sigma_0 \sigma_x] \asymp \exp(-|x|/\eta)$, with $\eta \equiv \eta(\nu) \to \infty$ as $\nu \to \nu_c$.
when $\nu = \nu_c$: $\mathbb{E}[\sigma_0 \sigma_x] \asymp |x|^{-2\delta}$ with $\delta = 1/8$.

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• Limit of geometric observables (connection probabilities, interfaces, \ldots) \sim SLE

A very simplified history:

• Lenz 1920 (the model) and Ising 1925 (one-dimensional case: no phase transition)

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- 1980'-1990': development of the CFT and its predictions of the scaling limits of the Ising correlation functions.
- 2000+: proof of the convergence of correlation functions to the predicted limits. prediction and proof of the scaling limits of the interfaces (SLE, CLE)

A *(finite) planar map* is a *proper embedding* of a finite connected graph into the sphere S_2 , seen up to the orientation-preserving homeomorphisms of S_2 .

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To avoid symmetry problems, we mark a corner (called the *root*) of the planar map. The resulting object is a *rooted planar map*, which will be called *map* in the sequel.

Image: A math a math

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A (rooted planar) *triangulation* is a map whose all faces are triangles.* When only the *internal* faces are triangles, we talk about *triangulation with boundary*.

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A planar triangulation



A quadrangulation with a boundary



A more general map with a boundary

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A planar triangulation





A quadrangulation with a boundary

A more general map with a boundary

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A *uniform triangulation* of size *n* is a random triangulation chosen uniformly among the (rooted planar) triangulations with *n* faces. A *Boltzmann triangulation* of weight *t* is a random triangulation chosen among all the triangulations with a probability $\mathbb{P}(\mathfrak{t}) \propto t^{\#faces(\mathfrak{t})}$.



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The partition function: $\mathcal{Z}(t) := \sum_{\mathfrak{t}} t^{\# faces(\mathfrak{t})} = \sum_{n} T_n \cdot t^n$, where $T_n := \# \{ \text{triangulations with } n \text{ faces} \}.$

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Asymptotic counting formula:

$$T_n \underset{n \to \infty}{\sim} c \cdot t_c^{-n} \cdot n^{-5/2}$$
 with $c = \frac{\sqrt{6}}{32\sqrt{\pi}}$ and $t_c = \frac{27}{256}$.

This is the standard asymptotic behavior of the number of planar maps:

- planar maps with *n* edges: $c = \frac{2}{\sqrt{\pi}}$ and $t_c = \frac{1}{12}$,
- triangulations with a perimeter of p and n internal faces: c = c(p) and $t_c = t_c(p)$,

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To be compared with the standard asymptotic behavior of the number of plane trees:

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 \sim *Two universality classes.* --> Other universal observables?

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Asymptotic counting formula:

$$T_n \sim c \cdot t_c^{-n} \cdot n^{-5/2}$$
 with $c = \frac{\sqrt{6}}{32\sqrt{\pi}}$ and $t_c = \frac{27}{256}$.

This is the standard asymptotic behavior of the number of planar maps:

- planar maps with *n* edges: $c = \frac{2}{\sqrt{\pi}}$ and $t_c = \frac{1}{12}$,
- triangulations with a perimeter of p and n internal faces: c = c(p) and $t_c = t_c(p)$,

To be compared with the standard asymptotic behavior of the number of *plane trees*:

#{plane trees with *n* edges}
$$\underset{n \to \infty}{\sim} c \cdot t_c^{-n} \cdot n^{-3/2}$$

 \sim Two universality classes. --> Other universal observables? \sim Yes, many!

Theorem (Aldous 1990, etc.)

Let \mathfrak{T}_n be a random tree of size *n* chosen uniformly from one of many classes of rooted plane trees,

$$\left(\mathfrak{T}_n, rac{C}{n^{1/2}} d_{g^r}^{\mathfrak{T}_n}\right) \stackrel{GM}{\underset{n \to \infty}{\longrightarrow}} (\mathcal{T}, d)$$

in distribution, where (\mathcal{T}, d) is Adlous' Continuous Random Tree (a compact metric space of Haudorff dimension 2, independent of the choice of the class).

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Theorem (Le Gall & Paulin 2008, Le Gall 2013, Miermont 2013, etc.)

Let \mathfrak{m}_n be a random map of size *n* chosen uniformly from one of many classes of rooted planar maps,

$$\left(\mathfrak{m}_n, rac{C}{n^{1/4}} d_{gr}^{\mathfrak{m}_n}
ight) \stackrel{GM}{\longrightarrow} (\mathcal{M}, \mathcal{D})$$

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A very simplified history:

1960': enumeration of maps using combinatorial methods (Tutte et al.)

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 \sim Yes. One way to do so: couple the law of the random map to a model of statistical physics, i.e. instead of choosing the random map *uniformly* or with a weight that only depends on its *size*, we choose $\mathbb{P}(\mathfrak{m}) \propto \mathcal{Z}_{\mathfrak{m}}$, where $\mathcal{Z}_{\mathfrak{m}}$ is the *partition function* of some statistical physics model living on the map \mathfrak{m} .

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 ~2000+: rigorous and more systematic methods for the enumeration of various classes of maps with additional structures.

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The model: Boltzmann Ising-triangulation

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Let \mathcal{T}_2 be the set of triangulations with a *simple* boundary, endowed with a partition of its boundary into 2 intervals. Denote by $p(\mathfrak{t})$ and $q(\mathfrak{t})$ the lengths of these intervals.



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Consider the set of Ising-decorated triangulations (Ising-triangulation for short):

$$\mathcal{IT}_{+-} = \{(\mathfrak{t}, \sigma) \mid \mathfrak{t} \in \mathcal{T}_2 \text{ and } \sigma \in \{+, -\}^{faces(\mathfrak{t})}\}$$

The elements of \mathcal{IT}_{+-} are endowed with *Dobrushin boundary condition*: we assign a sequence of p(t) spins + followed by q(t) spins - to the *outside of the boundary*. Let $\mathcal{E}(t, \sigma)$ be the number of monochromatic edges in (t, σ) .

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Let \mathcal{T}_2 be the set of triangulations with a *simple* boundary, endowed with a partition of its boundary into 2 intervals. Denote by $p(\mathfrak{t})$ and $q(\mathfrak{t})$ the lengths of these intervals.



Example:
$$(\mathfrak{t}, \sigma) \in \mathcal{IT}_{+-}$$

 $p(\mathfrak{t}) = 2, \ q(\mathfrak{t}) = 3,$
 $\# faces(\mathfrak{t}) = 7, \ \mathcal{E}(\mathfrak{t}, \sigma) = 5.$

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Boltzmann-Ising triangulation

Generating functions

$$Z(u, v; t, \nu) := \sum_{(\mathfrak{t}, \sigma) \in \mathcal{IT}_{\leftarrow}} u^{p(\mathfrak{t})} v^{q(\mathfrak{t})} t^{\#F(\mathfrak{t})} \nu^{\mathcal{E}(\mathfrak{t}, \sigma)}$$
$$Z_q(u; t, \nu) := [v^q] Z(u, v; t, \nu)$$
$$z_{p,q}(t, \nu) := [u^p v^q] Z(u, v; t, \nu) = [u^p] Z_q(u; t, \nu)$$

By convention $z_{0,0}(t, \nu) = Z(0, 0; t, \nu) = 1$.

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By convention $z_{0,0}(t, \nu) = Z(0, 0; t, \nu) = 1$.

For all $p, q \ge 0$ and $t, \nu > 0$ such that $z_{p,q}(t, \nu) < \infty$, we define a probability measure on the set $\{(\mathfrak{t}, \sigma) \in \mathcal{IT}_{+} \mid p(\mathfrak{t}) = p \text{ and } q(\mathfrak{t}) = q\}$ by

$$\mathbb{P}_{p,q}^{t,\nu}(\mathfrak{t},\sigma) = \frac{t^{\#F(\mathfrak{t})}\nu^{\mathcal{E}(\mathfrak{t},\sigma)}}{z_{p,q}(t,\nu)} \,.$$

We call a random variable of law $\mathbb{P}_{p,q}^{t,\nu}$ Boltzmann Ising-triangulation of (p,q)-gon.

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Previous results

Let $t_c(\nu)$ is the radius of convergence of $t \mapsto z_{l,0}(t,\nu)$. Let $\nu_c = 1 + 2\sqrt{7}$.

Theorem (Bernardi-Bousquet-Mélou 11, Albenque-Laurent-Schaeffer 18)

For all $\nu > 1$ and $(p,q) \neq (0,0)$, we have

$$\begin{bmatrix} t^n \end{bmatrix} z_{p,q}(t,\nu) \underset{n \to \infty}{\sim} \begin{cases} \kappa_{p,q}(\nu) \cdot t_c(\nu)^{-n} \cdot n^{-5/2} & (\nu \neq \nu_c) \\ \kappa_{p,q}(\nu_c) \cdot t_c(\nu_c)^{-n} \cdot n^{-7/3} & (\nu = \nu_c) \end{cases}$$

Moreover, $t_c(\nu)$ is C^2 -continuous on $(1, \infty)$ and analytic on $(1, \nu_c) \cup (\nu_c, \infty)$.

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(Easy) Corollaries

- $t_c(\nu)$ is also the radius of convergence of $t \mapsto z_{p,q}(t,\nu)$ and $z_{p,q}(t_c(\nu),\nu) < \infty$, for all $\nu > 1$ and $(p,q) \neq (0,0)$.
- $-\lim_{n\to\infty}\frac{1}{n}\log[t^n]z_{p,q}(t,\nu) = \log t_c(\nu)$. Thus $\log t_c(\nu)$ is the free energy density.

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We will focus on the case $t = t_c(\nu)$ ("maximal volume") and $\nu > 1$ (ferromagnetic). From now on, we assume $t = t_c(\nu)$ and omit the parameter t from the notations.

Main results

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Phase transition in the Ising model on a random 2D lattice

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Let $u_c(\nu)$ be the radius of convergence of $Z_0(u; \nu)$.

Continuous phase transition in the "surface tension" *

 $u_c(\nu)$ is positive and continuous on $(1,\infty)$, and is analytic everywhere except at ν_c .

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Theorem (Fixed-temperature asymptotics)

Fix $\nu > 1$. In the limit where $q \to \infty$ with p fixed, and then $p \to \infty$, we have

 $z_{p,q}(\nu) \sim a_p(\nu) \cdot u_c(\nu)^{-q} \cdot q^{-\alpha_0}$ and $a_p(\nu) \sim b(\nu) \cdot u_c(\nu)^{-p} \cdot p^{-\alpha_1}$.

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In the limit where $p, q \to \infty$ and $q/p \to \lambda$ for some fixed $\lambda \in (0, \infty)$, we have

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	$\nu < \nu_c$	$\nu = \nu_c$	$\nu > \nu_c$			
α_0	5/2	7/3	5/2		$(1+\lambda)^{-5/2}$	$(\nu < \nu_c)$
α_1	0	4/3	5/2	$c(\lambda) = \langle$	$\frac{4}{3}\int_0^\infty (1+r)^{-\frac{1}{3}} (\lambda+r)^{-\frac{1}{3}} dr$	$(\nu = \nu_c)$
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* In both of the two limits above, we have $-\lim_{\mu \to a} \log(z_{p,q}(\mu)) = \log u_{\mathfrak{g}}(\nu)$.

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Theorem (Near-critical asymptotics)

Fix $\beta \in (0,\infty]$ and assume $|\nu - \nu_c| = p^{-\beta}$. In the limit where $q \to \infty$ with p fixed and then $p \to \infty$, we have

$$z_{p,q}(\nu) \sim \tilde{a}_q(\beta) \cdot u_c(\nu)^{-p} \cdot p^{-lpha_0(eta)} \qquad and \qquad a_p(\nu) \sim \tilde{b}(\beta) \cdot u_c(\nu)^{-p} \cdot p^{- ilde{lpha}_1(eta)}$$

When $p, q \to \infty$ and $q/p \to \lambda$ for some fixed $\lambda \in (0, \infty)$, we have

$$z_{p,q}(\nu) \sim \tilde{c}(\lambda;\beta) \cdot u_c(\nu)^{-(p+q)} \cdot p^{-\alpha_2(\beta)}$$

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Local weak limit of the distributions $\mathbb{P}_{p,q}^{\nu}$

The *local distance* between two Ising-triangulations (\mathfrak{t}, σ) and (\mathfrak{t}', σ') is defined by

$$d_{\texttt{loc}}((\mathfrak{t},\sigma),(\mathfrak{t}',\sigma')) = 2^{-\sup\{r \in \mathbb{N} \colon B_r(\mathfrak{t},\sigma) = B_r(\mathfrak{t}',\sigma')\}}$$

where B_r denotes the ball of radius r (w.r.t. the graph distance) around the root.



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Local weak limit of the distributions $\mathbb{P}_{p,q}^{\nu}$

Theorem (Critical and off-critical local limits)

For each $\nu > 1$, one can construct probability distributions $(\mathbb{P}_p^{\nu})_{p\geq 0}$ and $\mathbb{P}_{\infty}^{\nu}$ such that $\mathbb{P}_{p,q}^{\nu} \xrightarrow[q \to \infty]{} \mathbb{P}_p^{\nu} \xrightarrow[p \to \infty]{} \mathbb{P}_{\infty}^{\nu}$ weakly with respect to the local distance. In the limit $p, q \to \infty$ and $q/p \to \lambda \in (0, \infty)$, the convergence becomes $\mathbb{P}_{p,q}^{\nu} \to \mathbb{P}_{\infty}^{\nu}$.

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- For $\nu > 1$ and $p < \infty$, \mathbb{P}_p^{ν} is supported on the set of one-ended triangulations with one infinite boundary (i.e. triangulations of the half plane).
- For $\nu \geq \nu_c$, $\mathbb{P}^{\nu}_{\infty}$ is also supported on the above set.
- For $\nu \in (1, \nu_c)$, $\mathbb{P}_{\infty}^{\nu}$ is supported on the set of two-ended triangulations.

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Theorem (Near-critical local limit)

When
$$\nu \to \nu_c$$
 at the same time as $p, q \to \infty$, we have $\mathbb{P}_{p,q}^{\nu} \xrightarrow[q \to \infty]{} \mathbb{P}_p^{\nu_c}$,
 $\mathbb{P}_p^{\nu} \xrightarrow[p \to \infty]{} \mathbb{P}_{\infty}^{\nu_c}$ and $\mathbb{P}_{p,q}^{\nu} \xrightarrow[p,q \to \infty]{} \mathbb{P}_{\infty}^{\nu_c}$ weakly with respect to the local distance

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Scaling limit of the main interface (work in progress)

Let $L_{p,q}^{\nu}$ be the length^{*} of the left-most Ising interface going from ρ to ρ' in a Boltzmann Ising-triangulation of law $\mathbb{P}_{p,q}^{\nu}$.

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Scaling limit of the main interface (work in progress)

Let $L_{p,q}^{\nu}$ be the length^{*} of the left-most Ising interface going from ρ to ρ' in a Boltzmann Ising-triangulation of law $\mathbb{P}_{p,q}^{\nu}$.

Theorem (off-critical and critical limit)

Fix $\nu > 1$ and $\lambda \in (0, \infty)$. In the limit $p, q \to \infty$ and $q/p \to \lambda$, the random variable $L_{p,q}^{\nu}/p$ converges in law to 0 if $\nu > \nu_c$, to a deterministic value $\ell(\lambda; \nu) > 0$ if $\nu < \nu_c$, and to the random variable of density $\frac{1}{\mathbb{Z}}(1+\mu x)^{-7/3}(\lambda+\mu x)^{-7/3}\mathbb{1}_{\{x>0\}}$ if $\nu = \nu_c$.

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Conjecture (near-critical limit)

 $\text{Fix }\beta>0 \text{ and }\lambda\in(0,\infty). \text{ In the limit } p,q\to\infty\text{, }q/p\to\lambda\text{ and }|\nu-\nu_c|=p^{-\beta}\text{,}$

- if $\nu > \nu_c$, or $\nu < \nu_c$ and $\beta > 1/3$, then $L_{p,q}^{\nu}/p^{\delta(\beta)}$ converges in distribution to a non-trivial random variable on $(0, \infty)$, where $\delta(\beta) = 2\alpha_0(\beta) \alpha_2(\beta) \in (0, 1]$.
- if $\nu < \nu_c$ and $\beta < \frac{1}{3}$, then $L_{p,q}^{\nu}/p$ converges to a deterministic value $\ell(\lambda; \beta) > 0$.

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Thank you for your attention !

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