# A Green-Kubo Formula for general Markov processes with continuous time parameter

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# 1 Background

### 1.1. What is Green-Kubo formula?

1827 R. Brown

"Pollen grains suspended in water perform a continual swarming motions."

1905, 1908, 1923, 1930, 1945 ... Einstein, Langevin, Smoluchowski, Wiener Ornstien, Uhlenbeck, Wang,



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x(t): the displacement of a Brownian particle of mass m is governed by Newton's second law

$$m\ddot{x} = -\frac{\dot{x}}{\mu} - \frac{\partial V}{\partial x} + f_{\text{Random}} \tag{1}$$

•  $-\frac{\dot{x}}{\mu}$ : Friction force due to viscosity of the fluid.

$$\mu = (6\pi\eta R)^{-1}$$
: mobility,  $\gamma = \frac{1}{m\mu}$ : frictional factor

- $-\frac{\partial V}{\partial x}$ : external potential V e.g. gravity
- ullet  $f_{
  m Random}$ : Random force of zero mean due to the impacts of fluid particles.



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## **Hypothesis:** $\gamma \gg 1$ (Overdamped)

The viscous term dominates the inertial one, so, ignore the acceleration.

Then, Eq. (1)

$$m\ddot{x} = -\frac{\dot{x}}{\mu} - \frac{\partial V}{\partial x} + f_{\text{Random}}$$

reduces to Langevin Equation

$$\dot{x} = -\mu \frac{\partial V}{\partial x} + \mu f_{\text{Random}}$$

$$= -\mu \frac{\partial V}{\partial x} + \eta(t)$$
(2)

Assume that  $\eta$ : zero mean and

$$\langle \eta(t), \eta(t') \rangle = D\delta(t - t')$$
 (3)

D is related to "Diffusion of the Brownian particle"



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If  $V(x) = x^2$ , (Particle connected to a Hookian spring), then, Eq. (2) is O-U equation

$$\dot{x} = -\mu x + \eta(t). \tag{4}$$

- Deterministic velocity:  $v(x) = -\mu x$
- Stochastic velocity  $\eta(t)$
- Mean position:  $\langle x(t)\rangle \equiv E(x(t)) = x(0)e^{-\mu t} \to 0 \text{ as } t \to \infty$
- Fluctuation around the mean position

$$\langle (x(t) - \langle x(t) \rangle)^2 \rangle \equiv Var(x(t)) \to \frac{D}{2\mu}, \text{ as } t \to \infty$$
 (5)



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In physics (canonical ensemble), once the Brownian particle reaches the equilibrium (stationary process) with the fluid at a temperature T, the probability of finding a Brownian particle at position x is given

$$\mathcal{P}_{\text{equilibrium}}(x) = \frac{1}{\sqrt{2\pi k_B T}} \exp\left[-\frac{x^2}{2k_B T}\right]$$

 $k_B$  is Boltzmann constant.

Then, 
$$\langle x^2 \rangle_{\text{equilibrium}} = k_B T$$
, by (5)  $\langle x^2 \rangle_{\text{equilibrium}} = \frac{D}{2\mu}$ 



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Now, obtaining the well-known Einstein relation

$$D = 2\mu k_B T \tag{6}$$

This equality connects the fluctuation of Noise (D) to the dissipation in the medium.

**Equilibrium**  $\Rightarrow$  **Balance** between **Fluctuation** and **Dissipative** 



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# Another derivation by Green (1952) and Kubo (1966) at the equilibrium state

• The mean velocity of x(t) at (t, x(t)) is

$$v(x(t)) \equiv -\mu x(t)$$

• The autocorrelation function of v(x(t)) is

$$R(t) \equiv E(v(x(0))v(x(t))) = \langle v(x(0))v(x(t))\rangle = \frac{\mu D^2}{2}e^{-\mu t}$$
 (7)

 $\Phi(\omega)$  is Fourier transform of R(t), it is the Power spectrum of v(x(t)).  $\Phi(\omega)$  can be regraded as "energy" in the frequency  $\omega$ .



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• moreover, by direct computation

$$R(t) = \langle v(x(0))^2 \rangle De^{-\mu t}$$
 (8)

• So,

$$\frac{D^2}{2} = \int_0^\infty R(t)dt = \langle v(x(0))^2 \rangle D\mu \tag{9}$$

- In fact,  $\langle v(x(0))^2 \rangle$  is kinetic energy, then  $\langle v(x(0))^2 \rangle = k_B T$
- Therefore, obtain Einstein relation  $D = 2\mu k_B T$  again.



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The formula

$$\frac{D^2}{2} = \int_0^\infty \langle v(0)v(t)\rangle dt = \int_0^\infty R(t)dt \tag{10}$$

is called **Green-Kubo formula** (Green-Kubo relation). It is a special form of **Fluctuation-Dissipative Theorem**.



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## 2 Main results

Stochastic processes and deterministic or random dynamical systems are considered as two of the most important mathematical approaches dealing with the problems in nonequilibrium statistical physics

- Qian, Guo, Guo, 1988, Green-Kubo formula holds for reversible diffusion processes
- Spohn, 1991, some linearized hydrodynamic system with fluctuation currents and hard core lattice gas in thermal equilibrium and
- Ruelle, 1999, smooth dynamical systems
- Qian, Qian, Zhang, 2003, ergodic **reversible**, **finite states** Markov Chains
- Jiang, Zhang, 2003, general **reversible** Markov processes by Spectral representation theory
- Chen, Chen and Qian, 2006, **irreversible** Markov Chains with **finite states** by Matrix analysis



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**Our main result:** A Green-Kubo formula holds for general Markov processes under some mild conditions.

**Problem:** How to understand "Mean Velocity" and "Diffusion Coefficient" of a general Markov process  $X_t$ ??

The state space is not  $\mathbf{R}^d$ .



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#### **Solution:**

We get information of Markov process  $X_t$  by observable

$$\varphi(X_t) = (\varphi_1, \varphi_2, \cdots, \varphi_m)(X_t)$$

Therefore, consider the "Mean Velocity" and "Diffusion Coefficient" of observable  $\varphi(X_t)$ .

Then, go back to Nelson 1966



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Let 
$$\mathcal{P}_t = \sigma\{X_s, s \leq t\}$$

## **Definition of Mean Forward Velocity of** $\varphi(X_t)$

$$V\varphi(X_t) = \lim_{\Delta t \downarrow 0} E\left[\frac{\varphi(X_{t+\Delta t}) - \varphi(X_t)}{\Delta t} \,|\, \mathcal{P}_t\right] \tag{11}$$

## **Definition of Conditional Diffusion Coefficient of** $\varphi(X_t)$

$$G\varphi(X_t) = \lim_{\Delta t \downarrow 0} \mathbb{E}\left[\frac{(\varphi(X_{t+\Delta t}) - \varphi(X_t))^T(\varphi(X_{t+\Delta t}) - \varphi(X_t))}{\Delta t} \middle| \mathcal{P}_t\right],\tag{12}$$



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 $X_t$ : stationary Markov process on a Polish space  $\Xi$ 

 $\mu$ : stationary probability distribution of  $X_t$ 

 $P_t$ : transition semigroup of  $X_t$  on  $L^2(\Xi, \mu)$ , assume it is strongly continuous.

A: generator of  $P_t$ .

 $\mathcal{D}(\mathcal{A})$ : the domain of the generator  $\mathcal{A}$ 



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# **Theorem: Green-Kubo formula**: For the stationary Markov process $\{X_t\}_{t\geq 0}$ , let

$$\mathcal{J} = \Big\{ f \in L^2(\Xi, \mu) : \lim_{t \to \infty} \|P_t f - \langle f \rangle_{\mu}\|_{L^2(\mu)} = 0 \Big\}.$$
 (13)

Then for any  $\varphi_i$ ,  $\varphi_j \in \mathcal{D}(\mathcal{A}) \cap \mathcal{J}$ , and  $\phi_i \phi_j \in \mathcal{D}(\mathcal{A})$ 

$$\langle G_{ij}\varphi(X_t)\rangle = \int_0^\infty \langle V\varphi_i(X_0)V\varphi_j(X_t)\rangle dt + \int_0^\infty \langle V\varphi_j(X_0)V\varphi_i(X_t)\rangle dt.$$
(14)

or,

$$\left( \langle G\varphi(X_t) \rangle \right)_{i,j=1\dots m} = \int_0^\infty \langle V\varphi(X_0) \otimes V\varphi(X_t) \rangle \, dt + \int_0^\infty \langle V\varphi(X_0) \otimes V\varphi(X_t) \rangle \, dt$$
(15)



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#### **Remark:**

• If  $X_t$  is reservable, Green-Kubo formula is

$$\langle G_{ij}\varphi(X_t)\rangle = 2\int_0^\infty \langle V\varphi_i(X_0)V\varphi_j(X_t)\rangle dt$$
 (16)

• The condition

$$\mathcal{J} = \left\{ f \in L^2(\Xi, \mu) : \lim_{t \to \infty} ||P_t f - \langle f \rangle_{\mu}||_{L^2(\mu)} = 0 \right\}$$
$$\varphi_i, \, \varphi_j \in \mathcal{D}(\mathcal{A}) \cap \mathcal{J}$$

is sharp.

For example,  $\dot{x}=1$ ,  $z(t)=e^{\mathrm{i}\,x(t)}$  can be regarded as a Markov process on  $\mathbb{S}^1$ . Lebesgue measure on  $\mathbb{S}^1$  is its stationary distribution. This example was given by Deuschel at this May for us.



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•  $X = \{X_t\}_{t \geq 0}$  is a stationary diffusion on  $\mathbb{R}^d$  with infinitesimal generator

$$\mathcal{A} = \frac{1}{2} \nabla \cdot (A(x) \nabla) + b(x) \cdot \nabla = \frac{1}{2} \sum_{i,j=1}^{d} \frac{\partial}{\partial x_i} a_{ij}(x) \frac{\partial}{\partial x_j} + \sum_{i=1}^{d} b_i(x) \frac{\partial}{\partial x_i},$$
(17)

$$\varphi_i(x) = x_i, \varphi_j(x) = x_j$$

Assume: the stationary distribution of  $X_t$  have a positive smooth density  $\rho(x)$ , and all the condition in Theorem are satisfied. Thus,



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$$V_i(x) \equiv V\varphi_i(x) = b_i(x) + \frac{1}{2} \sum_{k=1}^d \frac{\partial a_{ik}}{\partial x_k}(x), \tag{18}$$

$$V_j(x) \equiv V\varphi_j(x) = b_j(x) + \frac{1}{2} \sum_{k=1}^d \frac{\partial a_{jk}}{\partial x_k}(x), \tag{19}$$

$$G_{ij}\varphi(x) = a_{ij}(x) \tag{20}$$

#### Green-Kubo formula is

$$\int_{\mathbb{R}^d} a_{ij}(x)\rho(x)dx = \int_0^\infty \langle V_i(X_0)V_j(X_t)\rangle dt + \int_0^\infty \langle V_j(X_0)V_i(X_t)\rangle dt.$$
(21)



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#### **Sketch of Proof:**

• Due to direct computation

$$\langle G_{ij}\varphi(X_t)\rangle = -\langle \mathcal{A}\varphi_i, \varphi_j\rangle_{\mu} - \langle \mathcal{A}\varphi_j, \varphi_i\rangle_{\mu}.$$
 (22)

$$\int_{0}^{\infty} \langle V\varphi_{i}(X_{0})V\varphi_{j}(X_{t})\rangle dt = \int_{0}^{\infty} \langle P_{t}\mathcal{A}\varphi_{j}, \mathcal{A}\varphi_{i}\rangle_{\mu} dt (23)$$

$$= \lim_{t \to \infty} \langle P_{t}\varphi_{j}, \mathcal{A}\varphi_{i}\rangle_{\mu} - \langle \varphi_{j}, \mathcal{A}\varphi_{i}\rangle_{\mu}.$$
(24)

$$\lim_{t \to \infty} ||P_t f - \langle f \rangle_{\mu}||_{L^2(\mu)} = 0 \text{ AND } \langle \mathcal{A}\varphi \rangle_{\mu} = 0$$

$$\Rightarrow$$

$$\lim_{t \to \infty} \langle P_t f, \mathcal{A}\varphi \rangle_{\mu} = 0$$



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# 3 Examples

### 3.1. Stochastic evolution equations in Hilbert spaces

$$\begin{cases} dX_t = (AX_t + F(X_t)) dt + B dW_t, & t \ge 0, \\ X_0 = x \in H, \end{cases}$$
 (25)

**Reamrk:** Da Prato *et al*: Non symmetric dissipative stochastic systems.

Da Prato, G., Debussche A. and Goldys, B., *Some properties of invariant measures of non symmetric dissipative stochastic systems*. Probab. Theory Relat. Fields 123 (2002), 355-380.



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**Theorem: 3.1** (Da Prato, G., Debussche A. and Goldys, B.)

Under some conditions,.....  $\omega + \kappa < 0$  and  $C^{-1} \in L(H)$  and  $F \in C_b^2(H; H)$ , then we have for any  $f \in L^2(H, \mu)$ ,  $\mu$  is the invariant measure of  $X_t$ 

$$\lim_{t \to \infty} ||P_t f - \langle f \rangle_{\mu}||_{L^2(\mu)} = 0, \tag{26}$$

with exponential rate.

\* That is to say, it is the same as  $\mathcal{J}=L^2(H,\mu)$ . There are many observables f enough to guarantee Green-Kubo formula holds for Eq. (25)



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**Proposition 3.2** (Da Prato 2004) Assume that  $C \equiv BB^* = \mathbb{I}$  (identity operator) and  $F \in C_b(H, H)$ . Then

$$\mathcal{A}$$
 is symmetric if and only if  $F = \frac{1}{2}\nabla \log \rho$ ,

where  $\rho = \frac{d\mu}{d\nu}$ , the Radon-Nikodym derivative of  $\mu$  with respect to  $\nu$ ,  $\nu = N(0, Q_{\infty})$ .

**Example :**(Da Prato 2002) Stochastic Reaction-Diffusion Equation:  $H = L^2(0, 1)$ ,

$$\begin{cases}
Af = \frac{\partial^2}{\partial \zeta^2}(af) + b\frac{\partial}{\partial \zeta}f + cf, & f \in D(A) \\
D(A) = H^2(0,1) \cap H_0^1(0,1),
\end{cases}$$
(27)

•••

$$F(x) = \sum_{k=0}^{2m-1} a_k x^k, \ x \in \mathbb{R}, \ a_{2m-1} < 0$$



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## **3.2.** Critical interacting diffusion processes on $\mathbb{Z}^d$

Assume  $d \geq 3$ . Let  $X_t = \{X_t(i), i \in \mathbb{Z}^d\}$ 

$$dX_t(i) = \sum_{j \in \mathbb{Z}^d} q_{ij} X_t(i) dt + \sqrt{a(X_t(i))} dB_t(i), \qquad i \in \mathbb{Z}^d, \quad (28)$$

(H1)  $q_{ij} \ge 0$ ,  $i \ne j$ , and  $Q = \{q_{ij}\}$  is of finite range.

(H2) 
$$-q_{ii} = \sum_{j \neq i} q_{ij} = 1.$$

(H3) 
$$q_{ij} = q_{0(j-i)}$$
.

Shiga, T., An interacting system in population genetics. J. Math. Kyoto Univ. 20 (1980), 213-242.

Shiga, T., Shimizu, A., *Infinite dimensional stochastic differential equations and their applications*. J. Math. Kyoto Univ. 20 (1980), 395-416.

Deuschel, J. D., Invariance princible and empirical mean large deviation of the critical Ornstein-Uhlenbeck process. Ann. Probab. 17 (1989), 74-90.



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• The critical Ornstein-Uhlenbeck process: a(y) = 1

Proposition 3.3 (Chojnowska-Michalik, A., Goldys, B. 2002)

The following conditions are equivalent.

- (i) The semigroup  $P_t$  is symmetric in  $L^2(H, \mu)$ .
- (ii) If  $x \in \mathcal{D}(A^*)$  then  $Cx \in \mathcal{D}(A)$  and

$$ACx = CA^*x.$$

(iii) 
$$e^{tA}C = C(e^{tA})^*$$
 for all  $t \ge 0$ .

Chojnowska-Michalik, A., Goldys, B., *Symmetric Ornstein-Uhlenbeck semigroups and their generators*. Probab. Theory Relat. Fields 124 (2002), 459-486.

This is to say, there are many critical Ornstein-Uhlenbeck processes on  $\mathbb{Z}^d$  which are irreversible.



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• *The continuous time stepping stone model:* 

$$a(y) = y(1 - y), y \in [0, 1]$$

Feng, S., Schmuland, B., Vaillancourt, J., Zhou, X., *Reversibility of Interacting Fleming-Viot Processes with Mutation, Selection, and Recombination*. Arxiv preprint arXiv:0803.1492, 2008.

Their results imply this process is irreversible.



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• The measure-valued critical branching random walk: a(y) = yWe guess this process is irreversible, but we can not give a proof.

Deuschel (1989) show for all of the three above processes, for any  $f \in L^2(\Xi, \mu)$ ,  $\mu$  is the ergodic invariant measure.

$$\lim_{t\to\infty} ||P_t f - \langle f \rangle_{\mu}||_{L^2(\mu)} = 0.$$

That is to say, it is the same as  $\mathcal{J}=L^2(\Xi,\mu)$ . There are many observables f enough to guarantee Green-Kubo formula holds for for all of the three above processes



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## 4 Discussion

There are different understand for "Diffusion coefficient" and "Velocity" in statistical physics. So, different forms of Green-Kubo formulas are derived from different definitions.

Suppose that  $(v_x(t), v_y(t))$  is the velocity. Let the instantaneous position be

$$x(t) = x(0) + \int_0^t v_x(s)ds \equiv x(0) + \delta x(t).$$
 (29)

y(t) is similar.

**Definition:** The asymptotic diffusion coefficient is defined by

$$\lim_{t \to \infty} \frac{d}{dt} \langle \delta x^2(t) \rangle.$$



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### • Example 1 For the A-Langevin equations

$$\begin{bmatrix} d v_x(t) \\ d v_y(t) \end{bmatrix} = \begin{bmatrix} -r & \Omega \\ -\Omega & -r \end{bmatrix} \begin{bmatrix} v_x(t) \\ v_y(t) \end{bmatrix} dt + \sigma \begin{bmatrix} d B_1(t) \\ d B_2(t) \end{bmatrix}, \tag{30}$$

 $r>0, \Omega>0$ , Clearly,  $(v_x(t),v_y(t))$  has an invariant distribution. It is shown by Balescu that

$$\lim_{t \to \infty} \frac{d}{dt} \langle \delta x^2(t) \rangle = 2 \int_0^\infty \langle v_x(t) v_x(0) \rangle dt.$$
 (31)

This is also named Green-Kubo formula.

Balescu, R., *Statistical Dynamics Matter out of Equilibrium*. London: Imperial College Press, 1997.



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However, the process associated with a complete set of variables

$$\{\xi_t = (x(t), y(t), v_x(t), v_y(t))\}$$

is a Markov process. The mean forward velocity of x(t) is exactly  $v_x(t)$ . But, the process  $(\xi_t)_{t\geq 0}$  has no invariant distribution, which prevents our Green-Kubo formula from being used.



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 $\bullet$  **Example 2** Suppose p(t) is the momentum. For the harmonic oscillator, the explicit equations of motions are

$$\begin{cases} d x = \frac{p}{m} dt \\ d p = (-m\omega_0^2 x - \gamma p) dt + \sqrt{2kTm\gamma} dB(t), \end{cases}$$
 (32)

 $(\xi_t = (x(t), p(t)))_{t\geqslant 0}$  is a Markov diffusion with invariant distribution

$$\rho(x, p) = \frac{\omega_0}{2\pi kT} \exp\{-\frac{m\omega_0^2 x^2}{2kT} - \frac{p^2}{2mkT}\}\$$

Risken Risken, H., The Fokker-Plank Equation. 2nd Edition, Springer-Verlag, 1989.

Zwanzig, R., Nonequilibrium Statistical Mechanics. Oxford University Press, 2001.



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Then

$$V(x(t)) = \frac{p(t)}{m}$$

Thus, we have that

$$\langle Gx(t)\rangle = \lim_{t \to \infty} \frac{d}{dt} \langle \delta x^2(t)\rangle = 2 \int_0^\infty \langle \frac{p(t)}{m} \frac{p(0)}{m} \rangle dt.$$
 (33)

The asymptotic diffusion coefficient and the expectation of the conditional diffusion coefficient are identical now. This means these two different diffusion coefficients lead to the same Green-Kubo formula for this example.



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