# A Lévy-Fokker-Planck equation: entropies and convergence to equilibrium 

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Workshop on Markov Processes and Related Topics
Beijing Normal University, July 14-18, 2007
Work written in collaboration with C. Imbert from Paris-Dauphine

## Outline

(9) Introduction

- Ornstein-Uhlenbeck and Fokker-Planck equations
- Tools for the asymptotic behaviour
(2) The Lévy-Fokker-Planck equation
- The Lévy-Fokker-Planck equation
(3) Results
- The equilibrium
- Entropies
- Convergence towards the equilibrium


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## Ornstein-Uhlenbeck and Fokker-Planck equations

$$
\left\{\begin{array}{l}
d X_{t}=2 d B_{t}-\nabla V\left(X_{t}\right) d t \\
X_{0}=x
\end{array}\right.
$$

where $B_{t}$ is a standard Brownian Motion in $\mathbb{R}^{n}$.
Ito formula implies : the Semigroup $P_{t} f(x)=E_{x}\left(f\left(X_{t}\right)\right)$ satisfies the PDE

$$
\left\{\begin{array}{l}
\frac{\partial}{\partial \partial_{t}} P_{t} f(x)=L P_{t} f(x) \\
P_{0} f=f,
\end{array}\right.
$$

where $L f=\Delta f-\nabla V \cdot \nabla f$ is the IG of $P_{t}$. This is the Ornstein-Uhlenbeck equation.
Consider $L^{*}$ or $P_{t}^{*}$, the dual with respect to $d x$,

$$
\int L f g d x=\int f L^{*} g d x, \text { or } \int \mathrm{P}_{\mathrm{t}} \mathrm{fgdx}=\int \mathrm{fP}_{\mathrm{t}}^{*} \mathrm{gdx},
$$

then

$$
L^{*} g=\Delta g+\operatorname{div}(g . \nabla V)
$$

The Semigroup $P_{t}^{*} f(x)$ satisfies the PDE

$$
\left\{\begin{array}{l}
\frac{\partial}{\partial t} P_{t}^{*} f(x)=L^{*} P_{t}^{*} f(x) \\
P_{0}^{*} f=f
\end{array}\right.
$$

This is the Fokker-Planck equation.
Let $\mu_{V}=e^{-V} d x$ (assume that $\mu_{V}$ is a probability measure), $\left(P_{t}\right)_{t \geq 0}$ or $L$ is self adjoint in $L^{2}(d \mu v)$ and the by integration by parts

$$
\int L f g d \mu v=-\int \nabla f \cdot \nabla g d \mu V
$$

Under smooth assumptions :

$$
\lim _{t \rightarrow \infty} P_{t} f(x)=\int f d \mu_{V}
$$

or equivalently

$$
\lim _{t \rightarrow \infty} e^{V(x)} P_{t}^{*} g(x)=\int g d x
$$

The good question is HOW FAST?

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## Tools for the asymptotic behaviour

- Poincaré inequality : a $L^{2}$ convergence.

$$
\frac{d}{d t} \operatorname{var}_{\mu_{V}}\left(P_{t} f\right)=2 \int P_{t} f L P_{t} f d \mu_{V}-0=-2 \int\left|\nabla P_{t} f\right|^{2} d \mu_{V}
$$

If Poincaré inequality holds

$$
\begin{gathered}
\operatorname{var}_{\mu_{V}}(f) \leq C \int|\nabla f|^{2} d \mu_{V} \\
\operatorname{var}_{\mu_{V}}\left(P_{t} f\right) \leq e^{-2 t / C} \operatorname{var}_{\mu_{V}}(f)
\end{gathered}
$$

- Logarithmic Sobolev inequality a $L \log L$ convergence

$$
\frac{d}{d t} E n t_{\mu_{V}}\left(P_{t} f\right):=\frac{d}{d t} \int P_{t} f \log \frac{P_{t} f}{\int P_{t} f d \mu_{V}} d \mu_{V}=-4 \int\left|\nabla \sqrt{P_{t} f}\right|^{2} d \mu_{V}
$$

If Logarithmic Sobolev inequality holds

$$
\begin{gathered}
E n t_{\mu_{V}}\left(f^{2}\right) \leq C \int|\nabla f|^{2} d \mu_{V} \\
E_{n_{\mu_{V}}}\left(P_{t} f\right) \leq e^{-4 t / C} E_{\mu_{\nu V}}(f)
\end{gathered}
$$

When do we have a Poincaré or a logarithmic Sobolev inequality?

The well known Bakry-Emery $\Gamma_{2}$-criterion implies that if

$$
\operatorname{Hess}(\mathrm{V}) \geq \lambda \mathrm{Id},
$$

with $\lambda>0$ then logarithmic Sobolev inequality holds with $C=2 / \lambda$ and Poincaré inequality holds with $C=1 / \lambda$.

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## Definition of Lévy process

Lévy process $L_{t}=$ process with stationary \& indep increments
Fourier transform $\left(L_{t}\right)=e^{t \psi(\xi)}$ where $\psi$ is characterized by the Lévy-Khinchine formula.

$$
\psi(\xi)=-\sigma \xi \cdot \xi+i b \cdot \xi+\int\left(e^{i z \cdot \xi}-1-i z \cdot \xi \mathbf{1}_{B}(\xi)\right) \nu(d z)
$$

where $\nu$ is a singular measure satisfying

$$
\int_{B}|z|^{2} \nu(d z)<+\infty \quad \int_{\mathbb{R}^{d} \backslash B} \nu(d z)<+\infty
$$

$\sigma$ is a positive definite matrix and $b$ is a vector.
Parameters $(\sigma, b, \nu)$ characterize the Lévy process (or a inifinite divisible law).

- For all $t>0$ the law of $L_{t}$ is an infinite divisible law:
$\mu=\underbrace{\mu_{n} \star \cdots \star \mu_{n}}_{n \text { times }}$.
- Associated infinitesimal generators as for the Brownian Motion.

$$
\begin{aligned}
I(u)= & \operatorname{div}(\sigma \nabla u)+b \cdot \nabla u \\
& +\int\left(u(x+z)-u(x)-\nabla u(x) \cdot z 1_{B}(z)\right) \nu(d z)
\end{aligned}
$$

These operators appear everywhere (mathematical finance, mechanics, fluids etc.)

- Example : $(\sigma, b, \nu)=\left(0,0, \frac{1}{|z|^{\alpha+\alpha}} d z\right)$, the $\alpha$ stable process. In that case $\psi(\xi)=|\xi|^{\alpha}$. The case $\alpha=2$ is the Brownian motion. In that case $I(f)=\Delta^{\alpha / 2}(f)$, the fractional Laplacian.


## The Lévy-Fokker-Planck equation

Replace $\Delta$ by I a IG of a Lévy process in the Fokker-Planck equation :

$$
\left\{\begin{array}{l}
\frac{\partial}{\partial t} u=I(u)+\operatorname{div}(\mathrm{ux}) \\
u(0, x)=f(x)
\end{array}\right.
$$

The goal of this talk is to understand the asymptotic behaviour.
Remark: We assume that $V=x^{2} / 2$.
Questions:

- Find a steady state as $e^{-V}$ as for the classical case $\Delta$.
- Find the asymptotic behaviour of the Lévy-Fokker-Planck equation (LFP).
- Find conditions to get an asymptotic behaviour using inequalities as Poincaré or logarithmic Sobolev.


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An equilibrium $u_{\infty} \stackrel{\text { def }}{=}$ a stationary solution of the LFP $u_{\infty}$ can be seen as an invariant measure $\mu_{V}$ in the case of the Laplacian.
Proposition (Existence of an equilibrium)
Assume that

$$
\int_{\mathbb{R}^{d} \backslash B} \ln |z| \nu(d z)<+\infty
$$

There then exists an positive equilibrium $u_{\infty}$ :

$$
I\left(u_{\infty}\right)+\operatorname{div}\left(u_{\infty} x\right)=0 .
$$

Moreover, $u_{\infty} d x$ is an infinite divisible law whose characteristic exponent $A$ is

$$
A(\xi)=\int_{0}^{1} \psi(s \xi) \frac{d s}{s} .
$$

Of course the condition is satisfied in the case of the $\alpha$-stable. In that case $u_{\infty}$ is the infinite divisible law of the Lévy process, A $=\psi / \lambda$.

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For $\phi: \mathbb{R}^{+} \rightarrow \mathbb{R}$ convex and smooth and $\mu$ a probability measure, consider the $\phi$-entropy

$$
E_{\mu}^{\phi}(f)=\int \phi(f) d \mu-\phi\left(\int f d \mu\right)
$$

## - Examples

For $\phi(x)=\frac{1}{2} x^{2}\left(E_{\mu}^{\phi}=\right.$ the variance $), D_{\phi}(a, b)=\frac{1}{2}(a-b)^{2}$

$$
F_{\mu}^{\phi}(v)=\frac{1}{2} \iint(v(x+z)-v(x))^{2} \nu(d z) \mu(d x)
$$

For $\phi(x)=x \ln x-x-1\left(E_{\mu}^{\phi}=\right.$ entropy $), D_{\phi}(a, b)=a \ln \frac{a}{b}+b-a$
This is natural interpolation between the variance and the Entropy.

Define also a Bregman distance

$$
D_{\phi}(a, b)=\phi(a)-\phi(b)-\phi^{\prime}(b)(a-b) \geq 0
$$

## Theorem

Let $\mu(d x)=u_{\infty}(x) d x$, $\nu$ the Lévy measure associated to I and consider $v(t, x)=\frac{u(t, x)}{u_{\infty}(x)}$, then

$$
\frac{d}{d t} E_{\mu}^{\phi}(v(t, \cdot))=-\iint D_{\phi}(v(x+z), v(x)) \nu(d z) \mu(d x)
$$

- Fisher information

$$
F_{\mu}^{\phi}(v)=\iint D_{\phi}(v(x+z), v(x)) \nu(d z) \mu(d x)
$$

Can be seen as a Dirichlet form with respect to the measure $u_{\infty}(x) d x$

The proof of the theorem comes from

- A related equation : the Lévy-Ornstein-Ulenbeck equation (LOU)

The function $v=u / u_{\infty}$ satisfies

$$
\partial_{t} v=\frac{1}{u_{\infty}}\left(I\left(u_{\infty} v\right)-I\left(u_{\infty}\right) v\right)+x \cdot \nabla v \stackrel{\text { def }}{=} L v
$$

Dual operator of $L$ wrt $\mu$

$$
\int w_{1}\left(L w_{2}\right) d \mu=\int\left(\check{l}\left(w_{1}\right)-x \cdot \nabla w_{1}\right) w_{2} d \mu
$$

where $\check{I}$ is $I$ with $\check{\nu}(d x)=\nu(-d x)$.
Recall that in the classical case $L$ is a self-adjoint operator with respect to $\mu$.

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## Convergence towards the equilibrium

## Theorem

We assume that $\nu_{l}$ has a density $N$ with respect to $d x$ and satisfies

$$
\int_{\mathbb{R}^{d} \backslash B} \ln |z| N(z) d z<+\infty
$$

If $N$ is even and satisfies,

$$
\forall z, \quad \int_{1}^{+\infty} N(s z) s^{d-1} d s \leq C N(z)
$$

then for any smooth convex function $\Phi$ one gets :

$$
\forall t \geq 0, \quad \operatorname{Ent}_{u_{\infty}}^{\phi}\left(\frac{u}{u_{\infty}}\right) \leq e^{-\frac{t}{c} \operatorname{Ent}_{u_{\infty}}^{\phi}\left(\frac{u_{0}}{u_{\infty}}\right) . . . ~ . ~}
$$

## Proof

$$
\frac{d}{d t} E_{\mu}^{\phi}(v(t))=-F_{\mu}^{\phi}(t)=-\iint D_{\phi}(v(x+z), v(x)) \nu(d z) \mu(d x)
$$

it is enough to compare $F_{\mu}^{\phi}$ with $E_{\mu}^{\phi}$.

- A functional inequality [Wu'00,Chafaï'04]

If $\mu$ is an infinite divisible law $\phi$ satisfies $\phi^{\prime \prime}>0$ and $1 / \phi^{\prime \prime}$ concave on $\mathbb{R}^{+}$
Then

$$
E_{\mu}^{\phi}(f) \leq \iint D_{\phi}(v(x+z), v(x)) \nu_{\mu}(d z) \mu(d x)
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If

$$
\nu_{\mu} \leq C \nu_{l}
$$

This is true for fractional Laplacians and for Lévy process near $\alpha$-stable process.

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Then

$$
E_{\mu}^{\phi}(f) \leq C F_{\mu}^{\phi}(f)
$$

If

$$
\nu_{\mu} \leq C \nu_{l},
$$

then

$$
E_{\mu}^{\phi}\left(u / u_{\infty}\right) \leq E_{\mu}^{\phi}\left(\frac{u_{0}}{u_{\infty}}\right) e^{-\frac{t}{c}}
$$

This is true for fractional Laplacians and for Lévy process near $\alpha$-stable process.

