

Integration by parts formulae for the Wiener
measure restricted to domains in \mathbb{R}^d

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- $\mathcal{C} := C([0, 1]; \mathbb{R}^d)$.
- For $\Omega \subset \mathbb{R}^d$ an open region,

$$D := \{w \in \mathcal{C}; w(s) \in \Omega \text{ for all } 0 \leq s \leq 1\}.$$

1. IbP formula for the Wiener meas. on D

$$\left\{ \begin{array}{l} F : \text{a smooth functional on } \mathcal{C}. \\ h = (h_1, h_2, \dots, h_d), h_i \in C_0^\infty((0, 1)). \\ \text{For } a, b \in \Omega, \\ \mathcal{W}_{[0,1]}^{a,b} : \text{the pinned Wiener meas. s.t. } w(0) = a \text{ and } w(1) = b. \\ (\cdot, \cdot) : L^2([0, 1])\text{-inner product.} \end{array} \right.$$

$$\int_D \partial_h F(w) d\mathcal{W}_{[0,1]}^{a,b}(w) = - \int_D F(w) \sum_{i=1}^d (h_i'', w_i) d\mathcal{W}_{[0,1]}^{a,b}(w) + (\text{BT}),$$

where (BT) = Boundary Term

⇒ analogue to that in the Divergence Theorem
of finite dimension.

in the connection with:

Zambotti (PTRF 123, 2002):

$$\left\{ \Omega = (0, \infty) \subset \mathbb{R}, b = a \right.$$

$\left\{ \text{Biane's theorem relating Brownian and 3-dim. Bessel bridges} \right.$

⇒

SPDEs with reflection

of Nualart-Pardoux type

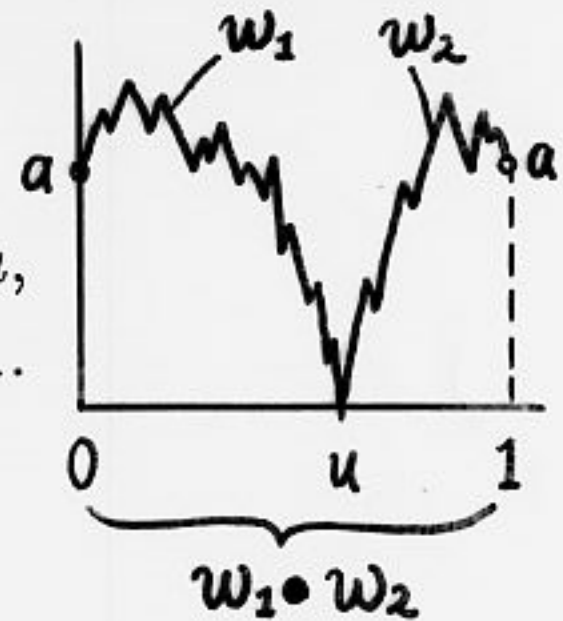
$$(1) \quad (BT) = - \int_0^1 du h(u) \frac{\sqrt{2}a^2}{\sqrt{\pi u^3(1-u)^3}} \exp\left(-\frac{a^2}{2u}\right) \exp\left\{-\frac{a^2}{2(1-u)}\right\} \\ \times E_{[0,u]}^{a,0,+} \otimes E_{[u,1]}^{0,a,+} [F(w_1 \bullet w_2)],$$

where

$E_{[l,r]}^{x,y,+}$: the expectation w.r.t. $P_{[l,r]}^{x,y,+}$, the law of 3-dim. Bessel bridge over $[l,r]$ s.t. $w(l) = x, w(r) = y$,

and

$$(2) \quad (w_1 \bullet w_2)(s) := \begin{cases} w_1(s), & 0 \leq s \leq u, \\ w_2(s), & u \leq s \leq 1. \end{cases}$$



An alternative expression of (1):

- B, \hat{B} : indep. Brownian motions starting at $a > 0$.
- Let

$$T_0(B) := \inf\{t > 0; B_t = 0\}, \quad T_0(\hat{B}) := \inf\{t > 0; \hat{B}_t = 0\}.$$

- Conditionally on $T_0(B) + T_0(\hat{B}) = 1$, define $Y = \{Y_t, 0 \leq t \leq 1\}$ by

$$(3) \quad Y_t := \begin{cases} B_t, & 0 \leq t \leq T_0(B), \\ \hat{B}_{T_0(B) + T_0(\hat{B}) - t}, & T_0(B) \leq t \leq T_0(B) + T_0(\hat{B}). \end{cases}$$



- $\mathbb{P}_{[0,1]}^{a,0,a}$: the law on $C([0,1]; [0,\infty))$ of Y .
- For a path $w \in C([0,1]; [0,\infty))$ starting at a and ending at b , and hitting the origin only one time, set

$S_0(w) :=$ the time at which $w(S_0(w)) = 0$

$\stackrel{\text{law}}{=} T_0(B)$ given $T_0(B) + T_0(\hat{B}) = 1$ (under $\mathbb{P}_{[0,1]}^{a,0,a}$).

- Conditionally on $T_0(B) + T_0(\hat{B}) = 1$, define $Y = \{Y_t, 0 \leq t \leq 1\}$ by

$$(3) \quad Y_t := \begin{cases} B_t, & 0 \leq t \leq T_0(B), \\ \hat{B}_{T_0(B) + T_0(\hat{B}) - t}, & T_0(B) \leq t \leq T_0(B) + T_0(\hat{B}). \end{cases}$$



- $\mathbb{P}_{[0,1]}^{a,0,a}$: the law on $C([0,1]; [0, \infty))$ of Y .
- For a path $w \in C([0,1]; [0, \infty))$ starting at a and ending at b , and hitting the origin only one time, set

$S_0(w) :=$ the time at which $w(S_0(w)) = 0$

$\stackrel{\text{law}}{=} T_0(B)$ given $T_0(B) + T_0(\hat{B}) = 1$ (under $\mathbb{P}_{[0,1]}^{a,0,a}$).

With these notations,

$$(4) \quad (1) = -4ae^{-2a^2} \mathbb{E}_{[0,1]}^{a,0,a} [h(S_0(w))F(w)].$$

This follows from the fact that, conditionally on $T_0(B) = u$, $\{B_t, 0 \leq t \leq T_0(B)\}$ has the same law as $P_{[0,u]}^{a,0,+}$, and that

$$P_a(T_0(B) \in du) = \frac{a}{\sqrt{2\pi u^3}} \exp\left(-\frac{a^2}{2u}\right) du.$$

Question: What about the case of more general Ω 's?

Answer: A similar expression to (4) holds.

Notations:

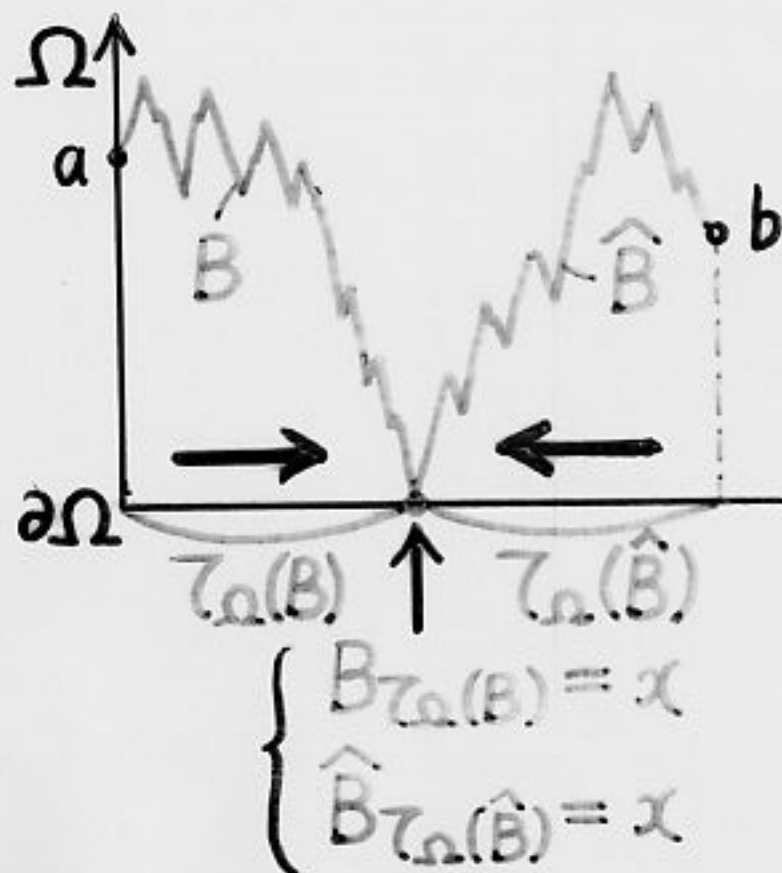
- $\Omega \subset \mathbb{R}^d$: an open, bounded region with a piecewise smooth boundary (so that the divergence theorem in finite dimension applies to Ω).
- B, \hat{B} : indep. d -dim. Brownian motions starting respectively at $a, b \in \Omega$.
- Let

$$\tau_{\Omega}(B) = \inf\{t > 0; B_t \in \mathbb{R}^d \setminus \Omega\}, \quad \tau_{\Omega}(\hat{B}) = \inf\{t > 0; \hat{B}_t \in \mathbb{R}^d \setminus \Omega\}.$$

- Given $\tau_{\Omega}(B) + \tau_{\Omega}(\hat{B}) = 1$, $B_{\tau_{\Omega}(B)} = x$, and $\hat{B}_{\tau_{\Omega}(\hat{B})} = x$, $Y = \{Y_t, 0 \leq t \leq 1\}$ is defined similarly to (3):

$$Y_t := \begin{cases} B_t, & 0 \leq t \leq \tau_{\Omega}(B), \\ \hat{B}_{\tau_{\Omega}(B) + \tau_{\Omega}(\hat{B}) - t}, & \tau_{\Omega}(B) \leq t \leq \tau_{\Omega}(B) + \tau_{\Omega}(\hat{B}). \end{cases}$$

- $\mathbb{P}_{[0,1]}^{a,x,b}$: the law on $C([0,1]; \bar{\Omega})$ of Y .



- For a path $w \in C([0, 1]; \overline{\Omega})$ starting at a and ending at b , and hitting the boundary $\partial\Omega$ only one time at the point x , define

$$S_x(w) := \text{the time at which } w(S_x(w)) = x.$$

- Δ_Ω : the Dirichlet Laplacian for Ω .
- $e^{-tH_\Omega}(y, z)$, $y, z \in \Omega$: the integral kernel of the semigroup e^{-tH_Ω} generated by $H_\Omega := -(1/2)\Delta_\Omega$; i.e.,

$$e^{-tH_\Omega}(y, z) = \frac{1}{\sqrt{(2\pi t)^d}} \exp\left\{-\frac{|y-z|^2}{2t}\right\} \mathcal{W}_{[0,t]}^{y,z} (w(s) \in \Omega, 0 \leq s \leq t).$$

Conditions on Ω (I still have to work on this part....):

- (i) for each fixed $t > 0$ and $y \in \Omega$, the integral kernel $e^{-tH_\Omega}(y, z)$ has an extension to $\bar{\Omega}$ which is C^1 up to the boundary;
- (ii) the restrictions to $\partial\Omega$ of functions which are harmonic on Ω , and C^1 up to boundary, is dense in the set of continuous functions on $\partial\Omega$.

- [Aizenman-Simon, Comm. Pure and Appl. Math. **35**]

Theorem A. 3. 2:

(i), (ii) \Rightarrow

$$(5) \quad P_a(\tau_\Omega(B) \in dt, B_{\tau_\Omega(B)} \in dx) = \frac{1}{2} \frac{\partial}{\partial \mathbf{n}_x} e^{-tH_\Omega}(a, x) \sigma(dx) dt,$$

where

$$\begin{cases} \sigma(dx) & : \text{ the surface measure on } \partial\Omega, \\ \mathbf{n}_x & : \text{ the inward normal vector at } x \in \partial\Omega, \\ \partial/\partial\mathbf{n}_x & : \text{ the normal derivative at } x. \end{cases}$$

Under these notations and conditions, (BT) is expressed as:

$$\begin{aligned} \text{(BT1)} \quad & -\frac{(2\pi)^{d/2}}{2} e^{|a-b|^2/2} \int_{\partial\Omega} \sigma(dx) \mathbb{E}_{[0,1]}^{a,x,b} [\mathbf{n}_x \cdot h(S_x(w)) F(w)] \\ & \times \int_0^1 du \frac{\partial}{\partial\mathbf{n}_x} e^{-uH_\Omega(a,x)} \frac{\partial}{\partial\mathbf{n}_x} e^{-(1-u)H_\Omega(b,x)}. \end{aligned}$$

2. Another expression of (BT) in terms of ground state for H_Ω

- e_0 : the smallest eigenvalue of $H_\Omega = -(1/2)\Delta_\Omega$; i.e.,

$$\begin{aligned} e_0 &= \inf \left\{ \frac{1}{2} \int_\Omega |\nabla f|^2 dx; f \in C_0^\infty(\Omega), \int_\Omega |f|^2 dx = 1 \right\} \\ &= - \lim_{t \rightarrow \infty} \frac{1}{t} \log P_a(\tau_\Omega(B) > t), \quad \forall a \in \Omega. \end{aligned}$$

- f_0 : a corresponding eigenfunction: $H_\Omega f_0 = e_0 f_0$.
- \mathbb{P}_x : the law on $C([0, \infty); \bar{\Omega})$ of X , the solution to the following SDE:

$$dX_t = dW_t + \frac{\nabla f_0}{f_0}(X_t) dt, \quad t \geq 0, \quad X_0 = x \in \bar{\Omega}.$$

- $p_t(x, y)$: the transition density of X .
- $\mathbb{P}_{[l,r]}^{x,y}$: the pinned measure over $[l, r]$ determined from \mathbb{P}_x so that

$$\mathbb{P}_{[l,r]}^{x,y}(w(l) = x, w(r) = y) = 1.$$

With these notations, (BT) is also expressed as:

$$\begin{aligned} \text{(BT2)} \quad & - \frac{(2\pi)^{d/2} e^{|a-b|^2/2 - e_0}}{2 f_0(a) f_0(b)} \int_0^1 du \int_{\partial\Omega} \sigma(dx) h(u) \cdot n_x \left(\frac{\partial f_0}{\partial n_x} \right)^2 \\ & \times \mathbb{E}_{[0,u]}^{a,x} \otimes \mathbb{E}_{[u,1]}^{x,b} [F(w_1 \bullet w_2)], \end{aligned}$$

where $w_1 \bullet w_2$ is defined as (2).

Example 1 (The case $\Omega = \{x \in \mathbb{R}; -1 < x < 1\}$).

- $f_0(x) = \cos(\frac{\pi}{2}x)$, $e_0 = \pi^2/8$.

- SDE:

$$dX_t = dW_t - \frac{\pi}{2} \tan(\frac{\pi}{2}X_t) dt, \quad t \geq 0, \quad X_0 = x \in [-1, 1].$$

$$\begin{aligned} \text{(BT)} = & - \left(\frac{\pi}{2}\right)^{5/2} e^{-\pi^2/8} \frac{e^{(a-b)^2/2}}{\cos(\frac{\pi}{2}a) \cos(\frac{\pi}{2}b)} \\ & \times \int_0^1 du h(u) \left\{ \mathbb{E}_{[0,u]}^{a,1} \otimes \mathbb{E}_{[u,1]}^{1,b} [F(w_1 \bullet w_2)] p_u(a, 1) p_{1-u}(b, 1) \right. \\ & \left. - \mathbb{E}_{[0,u]}^{a,-1} \otimes \mathbb{E}_{[u,1]}^{-1,b} [F(w_1 \bullet w_2)] p_u(a, -1) p_{1-u}(b, -1) \right\}. \end{aligned}$$

Derivation of (BT2) (heuristic):

"Penalty method" $\varphi_\varepsilon(x) := \frac{1}{\varepsilon} \mathbf{1}_{\Omega^c}(x)$ ($\varepsilon > 0$).

Consider the weight:

$$G_\varepsilon(w) := \exp\left\{-\int_0^1 \varphi_\varepsilon(w(s)) ds\right\}.$$

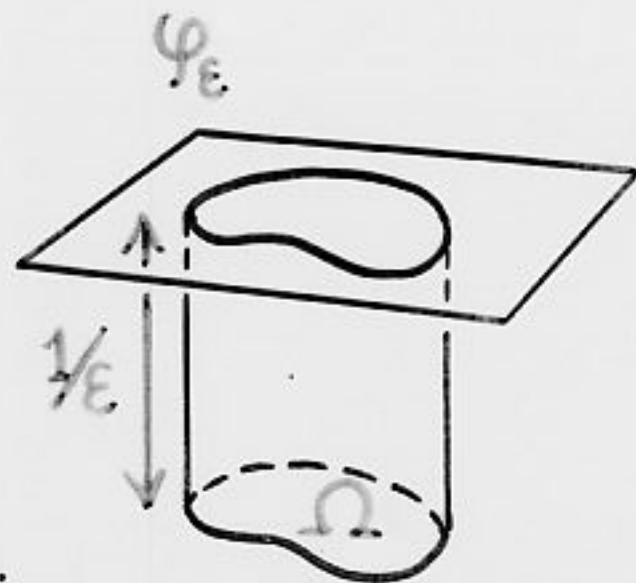
Cameron-Martin:

$$\xrightarrow{\varepsilon \downarrow 0} \mathbf{1}_D(w)$$

$$\int_{\mathcal{C}} F(w + \delta h) G_\varepsilon(w) d\mathcal{W}_{[0,1]}^{a,b}(w)$$

$$= \int_{\mathcal{C}} F(w) G_\varepsilon(w - \delta h)$$

$$\times \exp\left\{-\delta \sum_i (h_i'', w_i) - \frac{1}{2} \delta^2 \sum_i (h_i', h_i')\right\} d\mathcal{W}_{[0,1]}^{a,b}(w).$$



Take $\frac{d}{d\delta} \Big|_{\delta=0}$ on both sides:

RHS \rightarrow

$$\begin{aligned}
 & - \int_{\mathcal{C}} F(w) \sum_i (h_i'', w_i) G_\varepsilon(w) d\mathcal{W}_{[0,1]}^{a,b}(w) \\
 & + \sum_i \int_0^1 du h_i(u) \int_{\mathcal{C}} \partial_i \varphi_\varepsilon(w(u)) F(w) G_\varepsilon(w) d\mathcal{W}_{[0,1]}^{a,b}(w).
 \end{aligned}$$

$$\begin{aligned}
 \text{2nd term " = " } & - \int_0^1 du \int_{\partial\Omega} \sigma(dx) (h(u) \cdot \mathbf{n}_x) \frac{\mathcal{W}_{[0,1]}^{a,b}(w(u) \in dx)}{dx} \\
 & \times \frac{1}{\varepsilon} \mathcal{W}_{[0,u]}^{a,x} \otimes \mathcal{W}_{[u,1]}^{x,b} [F(w_1 \bullet w_2) G_\varepsilon(w_1 \bullet w_2)].
 \end{aligned}$$

(Can we confirm " = " by using, e.g., the theory of Watanabe's distribution?)

Rewrite:

$$\begin{aligned} & \frac{1}{\varepsilon} \mathcal{W}_{[0,u]}^{a,x} \otimes \mathcal{W}_{[u,1]}^{x,b} [F(w_1 \bullet w_2) G_\varepsilon(w_1 \bullet w_2)] \\ &= \frac{\mathcal{W}_{[0,u]}^{a,x} \otimes \mathcal{W}_{[u,1]}^{x,b} [F(w_1 \bullet w_2) G_\varepsilon(w_1 \bullet w_2)]}{\mathcal{W}_{[0,u]}^{a,x} \otimes \mathcal{W}_{[u,1]}^{x,b} [G_\varepsilon(w_1 \bullet w_2)]} \\ & \quad \times \frac{1}{\sqrt{\varepsilon}} \mathcal{W}_{[0,u]}^{a,x} [G_{\varepsilon,u}(w)] \times \frac{1}{\sqrt{\varepsilon}} \mathcal{W}_{[0,1-u]}^{x,b} [G_{\varepsilon,1-u}(w)] \\ &=: I_\varepsilon \times II_\varepsilon \times III_\varepsilon, \end{aligned}$$

where $G_{\varepsilon,t}(w) := \exp\{-\int_0^t \varphi_\varepsilon(w(s)) ds\}$, $0 < t < 1$.

Note: $I_\varepsilon \xrightarrow{\varepsilon \downarrow 0} \mathbb{E}_{[0,u]}^{a,x} \otimes \mathbb{E}_{[u,1]}^{x,b} [F(w_1 \bullet w_2)]$.

For II_ε and III_ε , we have:

Proposition 1. For $a \in \Omega$, $x \in \partial\Omega$, and $t > 0$,

$$\frac{1}{\sqrt{\varepsilon}} \mathcal{W}_{[0,t]}^{x,a} [G_{\varepsilon,t}(w)] \xrightarrow{\varepsilon \downarrow 0} \frac{p_t(x,a) e^{-e_0 t}}{\sqrt{2} f_0(a)} \frac{\partial f_0}{\partial \mathbf{n}_x} \times (2\pi t)^{d/2} e^{-|x-a|^2/2t}.$$

Note:

Proposition 1 \Rightarrow

$$II_\varepsilon \times III_\varepsilon \times \frac{\mathcal{W}_{[0,1]}^{a,b}(w(u) \in dx)}{dx} \xrightarrow{\varepsilon \downarrow 0} \frac{(2\pi)^{d/2}}{2} \frac{e^{-|a-b|^2/2 - e_0}}{f_0(a) f_0(b)} \left(\frac{\partial f_0}{\partial \mathbf{n}_x} \right)^2 p_u(a, x) p_{1-u}(b, x)$$

\Rightarrow (BT2) follows.

Key to Proposition 1:

$$H_\varepsilon := -\frac{1}{2}\Delta + \varphi_\varepsilon \quad \text{on} \quad L^2(\mathbb{R}^d; dx).$$

ε : sufficiently small so that H_ε has an ground state f_ε .

Lemma 2.

$$\frac{1}{\sqrt{\varepsilon}} f_\varepsilon(x) \xrightarrow{\varepsilon \downarrow 0} \frac{1}{\sqrt{2}} \frac{\partial f_0}{\partial \mathbf{n}_x} \quad \text{for } x \in \partial\Omega.$$

(f_0, f_ε : both normalized)

Remark 1. This may be regarded as a series expansion of f_ε in $\sqrt{\varepsilon}$:

$$f_\varepsilon(y) = f_0(y) + \sqrt{\varepsilon} (\text{1st factor}) + (\sqrt{\varepsilon})^2 (\text{2nd factor}) + \dots, \quad y \in \overline{\Omega}.$$

$$\text{Lemma 2} \iff (\text{1st factor})|_{\partial\Omega} = \frac{1}{\sqrt{2}} \frac{\partial f_0}{\partial \mathbf{n}_x}.$$

Is this kind of an expansion well-known? (WKB method?)

3. A rough sketch of the (computational) proof of (BT1)

Replace the time interval $[0, 1] \rightarrow [0, t], t > 0$.

Accordingly,

$$\mathcal{C} = C([0, t]; \mathbb{R}^d),$$

$$D = \{w \in C([0, t]; \mathbb{R}^d); w(s) \in \Omega, 0 \leq s \leq t\}.$$

In this case the IbP formula reads, for a smooth functional F on \mathcal{C} and $h = (h_1, \dots, h_d), h_i \in C_0^\infty((0, t))$,

$$\begin{aligned} \text{(IbP)} \quad & \int_D \partial_h F(w) d\mathcal{W}_{[0,t]}^{a,b}(w) \\ &= - \sum_{i=1}^d \int_0^t du \int_D h_i''(u) w_i(u) F(w) d\mathcal{W}_{[0,t]}^{a,b}(w) + (\text{BT}), \end{aligned}$$

$$\begin{aligned}
 (\text{BT}) = & -\frac{(2\pi t)^{d/2}}{2} \exp\left(\frac{|a-b|^2}{2t}\right) \int_{\partial\Omega} \sigma(dx) \mathbb{E}_{[0,t]}^{a,x,b}[\mathbf{n}_x \cdot h(S_x(w)) F(w)] \\
 & \times \int_0^t du \frac{\partial}{\partial \mathbf{n}_x} e^{-uH_\Omega}(a,x) \frac{\partial}{\partial \mathbf{n}_x} e^{-(t-u)H_\Omega}(b,x).
 \end{aligned}$$

Proof of (IbP):

- For a $\varphi : \mathbb{R}^d \rightarrow \mathbb{R}$ continuous, set $H_\Omega^\varphi := -(1/2)\Delta_\Omega + \varphi$.
- $e^{-tH_\Omega^\varphi}(y,z)$: the integral kernel of the semigroup $e^{-tH_\Omega^\varphi}, t > 0$:

$$\begin{aligned}
 e^{-tH_\Omega^\varphi}(y,z) = & \frac{1}{\sqrt{(2\pi t)^d}} \exp\left\{-\frac{|y-z|^2}{2t}\right\} \\
 & \times \mathcal{W}_{[0,t]}^{y,z} \left[e^{-\int_0^t \varphi(w(s)) ds}; w(s) \in \Omega, 0 \leq s \leq t \right].
 \end{aligned}$$

Lemma 3. For $a \in \Omega$ and $x \in \partial\Omega$,

$$E_a[e^{-\int_0^{\tau_{\Omega}(B)} \varphi(B_s) ds} | \tau_{\Omega}(B) = u, B_{\tau_{\Omega}(B)} = x] = \frac{\frac{\partial}{\partial \mathbf{n}_x} e^{-uH_{\Omega}^{\varphi}(a, x)}}{\frac{\partial}{\partial \mathbf{n}_x} e^{-uH_{\Omega}(a, x)}}.$$

Remark 2. By L'Hospital's rule, RHS is equal to

$$\lim_{\substack{z \rightarrow x \\ z \in \Omega}} \frac{\mathcal{W}_{[0, u]}^{a, z}[e^{-\int_0^u \varphi(w(s)) ds}; w(s) \in \Omega, 0 \leq s \leq u]}{\mathcal{W}_{[0, u]}^{a, z}(w(s) \in \Omega, 0 \leq s \leq u)}.$$

Proof. Let f : a test function on $\partial\Omega$. By (5),

$$\begin{aligned} & E_a[e^{-\int_0^{\tau_\Omega(B)} \varphi(B_s) ds} f(B_{\tau_\Omega(B)}); \tau_\Omega(B) > t] \\ &= \frac{1}{2} \int_t^\infty du \int_{\partial\Omega} \sigma(dx) f(x) \frac{\partial}{\partial \mathbf{n}_x} e^{-uH_\Omega}(a, x) \\ & \quad \times \mathbb{E}_a[e^{-\int_0^{\tau_\Omega(B)} \varphi(B_s) ds} | \tau_\Omega(B) = u, B_{\tau_\Omega(B)} = x]. \end{aligned}$$

So it suffices to prove

$$\begin{aligned} (6) \quad & E_a[e^{-\int_0^{\tau_\Omega(B)} \varphi(B_s) ds} f(B_{\tau_\Omega(B)}); \tau_\Omega(B) > t] \\ &= \frac{1}{2} \int_t^\infty du \int_{\partial\Omega} \sigma(dx) f(x) \frac{\partial}{\partial \mathbf{n}_x} e^{-uH_\Omega^\varphi}(a, x). \end{aligned}$$

Note: $\{\tau_\Omega(B) > t\} = \{B_s \in \Omega, 0 \leq s \leq t\}$. So,

LHS of (6)

$$\begin{aligned}
 &= E_a[U(B_t) e^{-\int_0^t \varphi(B_s) ds}; B_s \in \Omega, 0 \leq s \leq t] \quad (\because \text{Markov property}) \\
 &= \int_\Omega dx U(x) e^{-tH_\Omega^\varphi(a, x)}, \quad (\text{by definition})
 \end{aligned}$$

where $U(x) := E_x[f(B_{\tau_\Omega(B)}) e^{-\int_0^{\tau_\Omega(B)} \varphi(B_s) ds}]$. Thus

$$\begin{aligned}
 \frac{\partial}{\partial t}(\text{LHS of (6)}) &= \int_\Omega dx U(x) (-H_\Omega^\varphi) e^{-tH_\Omega^\varphi(a, x)} \\
 &= -\frac{1}{2} \int_{\partial\Omega} \sigma(dx) f(x) \frac{\partial}{\partial \mathbf{n}_x} e^{-tH_\Omega^\varphi(a, x)} \\
 &= \frac{\partial}{\partial t}(\text{RHS of (6)}).
 \end{aligned}$$

The 2nd equality follows from the divergence theorem and the fact that U solves: $H_{\Omega}^{\varphi} U = 0$ on Ω , $U|_{\partial\Omega} = f$. Q.E.D.

Back to the proof of (IbP):

Note: (IbP) is linear in F and $h = (h_i)_{1 \leq i \leq d}$.

↓

Approximate F and each h_i by linear combinations of the form:

$$F(w) \sim \sum_j e^{-\int_0^t \varphi_j(w(s)) ds}, \quad \varphi_j : \text{continuous},$$

$$h_i(u) \sim \sum_j k_i^j(u) l_i^j(t-u), \quad 0 \leq u \leq t, \quad k_i^j, l_i^j \in C_0^{\infty}((0, \infty)).$$

So it suffices to prove (IbP) for

$$F(w) = e^{-\int_0^t \varphi(w(s)) ds} \quad \text{and} \quad h_i(u) = k_i(u) l_i(t-u), \quad 1 \leq i \leq d :$$

$$(7) \quad \sum_{i=1}^d I_1^i(t) = \sum_{i=1}^d I_2^i(t) + \sum_{i=1}^d I_3^i(t),$$

where

$$I_1^i(t)$$

$$= (2\pi t)^{-\frac{d}{2}} e^{-\frac{|a-b|^2}{2t}} \int_0^t du k_i(u) l_i(t-u) \int_D \partial_i \varphi(w(u)) e^{-\int_0^t \varphi(w(s)) ds} d\mathcal{W}_{[0,t]}^{a,b}(w),$$

$$I_2^i(t)$$

$$= (2\pi t)^{-\frac{d}{2}} e^{-\frac{|a-b|^2}{2t}} \int_0^t du \{k_i''(u) l_i(t-u) + k_i(u) l_i''(t-u) - 2k_i'(u) l_i'(t-u)\} \\ \times \int_D w_i(u) e^{-\int_0^t \varphi(w(s)) ds} d\mathcal{W}_{[0,t]}^{a,b}(w),$$

and by Lemma 3,

$$I_3^i(t) = \frac{1}{2} \int_0^t du \int_{\partial\Omega} \sigma(dx) n_x^i k_i(u) l_i(t-u) \frac{\partial}{\partial n_x} e^{-uH_\Omega^\varphi}(a, x) \frac{\partial}{\partial n_x} e^{-(t-u)H_\Omega^\varphi}(b, x).$$

Here $\partial_i = \partial/\partial x_i$, n_x^i : the i -th component of n_x .

Note:

$$I_1^i(t) = \int_0^t du k_i(u) l_i(t-u) \int_\Omega dx \partial_i \varphi(x) e^{-uH_\Omega^\varphi}(a, x) e^{-(t-u)H_\Omega^\varphi}(b, x),$$

$$I_2^i(t) = (\text{similar expression in terms of the integral kernel}).$$

So, by taking the Laplace transform in t ,

(7) \iff

$$(8) \quad \sum_{i=1}^d II_1^i(\gamma) = \sum_{i=1}^d II_2^i(\gamma) + \sum_{i=1}^d II_3^i(\gamma), \quad \gamma > 0,$$

where

$$II_1^i(\gamma) = \int_0^\infty \int_0^\infty ds dt e^{-\gamma(s+t)} k_i(s) l_i(t) \int_\Omega dx \partial_i \varphi(x) e^{-sH_\Omega^\varphi(a, x)} e^{-tH_\Omega^\varphi(b, x)},$$

$$II_2^i(\gamma) = \int_0^\infty \int_0^\infty ds dt e^{-\gamma(s+t)} \{k_i''(s)l_i(t) + k_i(s)l_i''(t) - 2k_i'(s)l_i'(t)\} \\ \times \int_\Omega dx x_i e^{-sH_\Omega^\varphi(a, x)} e^{-tH_\Omega^\varphi(b, x)},$$

$$II_3^i(\gamma) = \frac{1}{2} \int_0^\infty \int_0^\infty ds dt e^{-\gamma(s+t)} k_i(s) l_i(t) \\ \times \int_{\partial\Omega} \sigma(dx) n_x^i \frac{\partial}{\partial n_x} e^{-sH_\Omega^\varphi(a, x)} \frac{\partial}{\partial n_x} e^{-tH_\Omega^\varphi(b, x)}.$$

Note: by simple integration by parts,

$$\begin{aligned}
 II_2^i(\gamma) = & \\
 & - \int_0^\infty \int_0^\infty ds dt e^{-\gamma(s+t)} k_i(s) l_i(t) \int_\Omega dx \left\{ (H_\Omega^\varphi e^{-sH_\Omega^\varphi}(a, x)) (\partial_i e^{-tH_\Omega^\varphi}(b, x)) \right. \\
 & \left. + (\partial_i e^{-sH_\Omega^\varphi}(a, x)) (H_\Omega^\varphi e^{-tH_\Omega^\varphi}(b, x)) \right\}.
 \end{aligned}$$

Compare these expressions \Rightarrow (8) is reduced to:

Proposition 4. For functions f, g on $\bar{\Omega}$ s.t. $f|_{\partial\Omega} = g|_{\partial\Omega} = 0$, and for $v \in \mathbb{R}^d$,

$$\begin{aligned}
 (9) \quad \int_\Omega dx (v \cdot \nabla \varphi) fg = & - \int_\Omega dx \left\{ (v \cdot \nabla g) H_\Omega^\varphi f + (v \cdot \nabla f) H_\Omega^\varphi g \right\} \\
 & + \frac{1}{2} \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) \frac{\partial f}{\partial \mathbf{n}_x} \frac{\partial g}{\partial \mathbf{n}_x}.
 \end{aligned}$$

Proof. We only need to consider the case $g = f$. Then by the definition of H_{Ω}^{φ} ,

$$\text{1st term on RHS} = \int_{\Omega} dx (v \cdot \nabla f) \Delta f - 2 \int_{\Omega} dx (v \cdot \nabla f) \varphi f.$$

Note: 2nd term = LHS of (9); indeed,

$$\begin{aligned} & \int_{\Omega} dx (v \cdot \nabla \varphi) |f|^2 + 2 \int_{\Omega} dx (v \cdot \nabla f) \varphi f \\ &= \int_{\Omega} dx \operatorname{div}(\varphi |f|^2 v) \\ &= - \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) \varphi |f|^2 && (\because \text{divergence theorem}) \\ &= 0. && (\because f|_{\partial\Omega} = 0) \end{aligned}$$

So It remains to prove:

$$(10) \quad \int_{\Omega} dx (v \cdot \nabla f) \Delta f + \frac{1}{2} \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) \left| \frac{\partial f}{\partial \mathbf{n}_x} \right|^2 = 0.$$

Apply the divergence theorem to $(v \cdot \nabla f) \nabla f$ to get:

(11)

$$\begin{aligned} & \int_{\Omega} dx (v \cdot \nabla f) \Delta f + \sum_{i,j=1}^d \int_{\Omega} v_i (\partial_i \partial_j f) \partial_j f \\ &= - \int_{\partial\Omega} \sigma(dx) (v \cdot \nabla f) \frac{\partial f}{\partial \mathbf{n}_x} \\ &= - \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) \left| \frac{\partial f}{\partial \mathbf{n}_x} \right|^2. \quad (\because \nabla f = \frac{\partial f}{\partial \mathbf{n}_x} \mathbf{n}_x \text{ at } x \in \partial\Omega) \end{aligned}$$

Moreover,

$$\begin{aligned} (12) \quad \sum_{i,j=1}^d \int_{\Omega} dx v_i (\partial_i \partial_j f) \partial_j f &= \frac{1}{2} \int_{\Omega} dx \operatorname{div}(|\nabla f|^2 v) \\ &= -\frac{1}{2} \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) |\nabla f|^2 \\ &= -\frac{1}{2} \int_{\partial\Omega} \sigma(dx) (v \cdot \mathbf{n}_x) \left| \frac{\partial f}{\partial \mathbf{n}_x} \right|^2. \end{aligned}$$

Combining (11) and (12) yields (10) and ends the proof of the proposition. Q.E.D.